# Modeling Photochemical [4 + 4] Cycloadditions: Conical Intersections Located with CASSCF for Butadiene + Butadiene 

Michael J. Bearpark, ${ }^{\dagger}$ Merce Deumal, ${ }^{\dagger}, \stackrel{+}{i}$ Michael A. Robb, ${ }^{*}, \dagger$ Thom Vreven, ${ }^{\dagger}$ Naoko Yamamoto, ${ }^{\dagger}$ Massimo Olivucci, $*, \S$ and Fernando Bernardi ${ }^{\S}$<br>Contribution from the Department of Chemistry, King's College London, Strand, London WC2R 2LS, U.K., and Dipartimento di Chimica 'G. Ciamician' dell' Universitá di Bologna, Via Selmi 2, 40126 Bologna, Italy

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#### Abstract

The $[4+4]$ photocycloaddition of butadiene + butadiene has been studied at the CASSCF /4-31G level, as a prototype for a class of photochromic systems. For this model system, minima and transition structures are characterized by analytic frequency calculations, and conical intersections are located. Our results indicate that the standard model for the $[4+4]$ addition (based on $\mathrm{H}_{4}$ ) needs to be revised. The reorganization of all $8 \pi$ electrons is crucial (i.e., it is not always the same $4 \pi$ electrons that are important). Efficient nonradiative decay of butadiene + butadiene can be explained by the presence of two distinct $S_{1} / S_{0}$ conical intersections. The first-the lowestenergy point on $S_{1}$ overall-is preceded by a barrier for the formation of a new $\sigma$ bond. The resulting structure is similar to those previously characterized for methyl migration in but-1-ene and the addition of ethylene to benzene. A higher-energy barrier leads to a second crossing which resembles the rhomboidal funnel for the [ $2+2$ ] addition of ethylene + ethylene, but which involves only one double bond from each butadiene. Both reaction paths commence at a true pericyclic minimum, at which the $\left(S_{0}-S_{1}\right)$ energy gap of $\sim 37 \mathrm{kcal} \mathrm{mol}^{-1}$ prohibits decay.


## Introduction

The $\left[4_{s}+4_{s}\right]$ cycloaddition is thermally forbidden but photochemically allowed. ${ }^{1,2}$ In this work we shall consider the prototypical $[4+4]$ cycloaddition of two butadienes which serves as a model for many examples of more complex [4+4] cycloadditions such as anthracene photodimerization.

In spite of the fact that one has a formal $8 \pi$ electron process, it is usually assumed that only $4 \pi$ electrons are actually involved in the recoupling process along the reaction path so that the reactivity can be rationalized in a similar fashion to the classic $[2+2]$ cycloaddition. The conventional model for the [4+ 4] photocycloaddition process ${ }^{3 \mathrm{a}, 2 \mathrm{a}, 4}$ is based on VB calculations for the $\mathrm{H}_{4}$ system $^{3}$ (Figure 1) which indicated the presence of a pericyclic minimum in the region of $S_{1}$ corresponding to a doubly-excited state. This minimum is assumed to be the decay funnel for the photochemical process. ${ }^{5}$ However, the original $\mathrm{H}_{4}$ calculations illustrate the problems with a symmetry constraint: the pericyclic minimum in Figure 1 is in fact a

[^0]
## Scheme 1


transition structure on a rhomboidal distortion coordinate ${ }^{3,6 a b}$ which connects two lower-energy conical intersections. ${ }^{7}$ These intersections have since been characterized for the ethylene + ethylene system (Scheme 1a) ${ }^{6 \mathrm{a}, \mathrm{b}}$ and provide an efficient mechanism for nonradiative decay to $S_{0}$.

The purpose of this paper is to examine the $[4+4]$ addition of butadiene + butadiene $^{8}$ as a model for general $[4+4]$ cycloadditions. This model is small enough to permit full geometry optimization at the CASSCF level for minima,

## Scheme 2



Scheme 3

transition structures, and conical intersections. ${ }^{6}$ We shall show that a model involving $4 \pi$ electron recoupling is indeed adequate to describe $[4+4]$ photochemical processes and that the interesting regions of the potential energy surfaces for real [4 $+4]$ systems will be determined by the constraints of the $\sigma$-bonded framework. However, the $4 \pi$ electrons are not in general the same as for a $[2+2]$ reaction, and generalization of this model ${ }^{4}$ is not simple. For the ethylene + ethylene system (Scheme 1a), the pericyclic structure is a transition structure on $S_{1}$ that connects two rhomboid conical intersections. In contrast we find a true pericyclic $\left[4_{s}+4_{s}\right]$ minimum $A$ for butadiene + butadiene (Scheme 1b). Rhomboidal distortion does not lead to a surface crossing in this case. Rather, two reaction paths exist, commencing at the pericyclic minimum $\mathbf{A}$ and proceeding via low barriers to the conical intersections $\mathbf{C}$ ( $\left.\left[22_{\mathrm{s}}+2_{\mathrm{s}}\right]\right)$ and $\mathbf{E}([4+4])$ illustrated in Scheme 2.

[^1]Scheme 4


The electronic origin of the crossings $\mathbf{C}$ and $\mathbf{E}$ (Scheme 2) can be related to four-electron recoupling processes in other systems we have studied before. ${ }^{6 a-c}$ Both $\mathbf{C}$ and $\mathbf{E}$ are tetraradicaloid, decay from which can lead to a mixture of products as there are several possibilities for electron recoupling on $S_{0}$. The conical intersection $\mathbf{E}$ (Scheme 2) is the lowestenergy point on $S_{1}$ overall. One $\sigma$ bond has already been formed, and the four radical centers are arranged in the manner of the conical intersection for methyl migration in but-1-ene shown at the bottom of Scheme 2. A similar crossing is found to be the global minimum of $S_{1}$ for the ethylene + benzene system. ${ }^{6 \mathrm{~m}}$ The higher-energy crossing $\mathbf{C}$ resembles the one due the rhomboidal distortion in the $\mathrm{H}_{4}$ system. ${ }^{3}$ One double bond participates from each butadiene molecule; the other spectates.

Little experimental information is available on the direct photochemical dimerization of butadiene + butadiene ${ }^{9,2 b, 10}$ (in contrast with the triplet sensitised reaction ${ }^{11}$ ). This is due in part to the rapid unimolecular decay of photoexcited butadiene itself: Mathies has established ${ }^{12}$ that crossing from the initially excited $1^{1} B_{u}$ state of butadiene to the reactive $2^{1} \mathrm{~A}_{\mathrm{g}}$ state takes place on the same time scale as vibrational relaxation ( $\sim 10 \mathrm{fs}$ ), ${ }^{6 e}$

[^2]Table 1. Energies of the Optimized CASSCF/4-31G Structures ${ }^{i}$

|  | figure | active space |  |  | $S_{0}, E_{\mathrm{h}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |

${ }^{a}$ Energy on $\mathrm{S}_{1}$ relative to pericyclic minimum A. ${ }^{b} \mathrm{M}=$ minimum; $\mathrm{TS}=$ transition structure; $\mathrm{X}=$ conical intersection. ${ }^{c}$ Separation $10 \AA .{ }^{d} \mathrm{~m}$ CI (n) indicates an m-orbital CASSCF calculation carried out with orbitals optimized for a n-orbital active space. ${ }^{e}$ Relative energies calculated with the CAS8 active space. ${ }^{f}$ tc $=$ trans + cis suprafacial approach. ${ }^{g}$ All structures optimized on $S_{1}$, unless indicated by $/ \mathrm{S}_{0} .{ }^{h}$ State-averaged orbitals. ${ }^{i}$ Analytic frequency calculations carried out for $\mathbf{4 a}, \mathbf{5 a}, \mathbf{5 b}$, and $\mathbf{8}$. The minimum $\mathbf{8}$ and transition structure $\mathbf{5 b}$ are almost coincident.

Table 2. CASSCF/6-31G* Energies at the Optimized CASSCF/4-31G Structures on $S_{1}$

| structure | figure | active space | $S_{0}, E_{\mathrm{h}}$ | $S_{1}, E_{\mathrm{h}}$ | $\begin{gathered} \Delta E\left(S_{1}-S_{0}\right) \\ \text { kcal } \mathrm{mol}^{-1} \end{gathered}$ | $\begin{aligned} & \Delta E \text { on } \mathrm{S}_{1}{ }^{a} \\ & \mathrm{kcal} \mathrm{~mol}^{-1} \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{M}^{b} \mathbf{A}$ | 4a | 8 |  | -309.78750 |  | 0.0 |
| $\infty^{c}$ | 4b | 8 |  | -309.76193 |  | +16.0 |
| TS B | 5 a | 8 |  | -309.76336 |  | +15.1 |
| TS D | 5b | 8 |  | -309.77691 |  | +6.6 |
| X C | 6 | 8 | -309.77776 ${ }^{f}$ | $-309.77067^{f}$ | 4.4 | +10.6 |
| X E | 6 b | $6$ | $-309.79459^{f}$ | $-309.79286^{f}$ | 1.1 |  |
|  |  | $8 \mathrm{CI}(6)^{d}$ | $-309.81067{ }^{f}$ | $-309.80890{ }^{f}$ | 1.1 | $-13.4{ }^{e}$ |
| X | 12 | 4 | $-309.69086^{f}$ | $-309.68443^{f}$ | 3.9 |  |
|  |  | $8 \mathrm{CI}(4)^{d}$ | -309.72120 ${ }^{\text {f }}$ | $-309.71453^{f}$ | 4.2 | +45.8 |

${ }^{a}$ Energy on $S_{1}$ relative to pericyclic minimum A. ${ }^{b} \mathrm{M}=$ minimum; $\mathrm{TS}=$ transition structure; $\mathrm{X}=$ conical intersection. ${ }^{c}$ Separation $10 \AA .{ }^{d} \mathrm{~m}$ CI (n) indicates an m-orbital CASSCF calculation carried out with orbitals previously optimized for a n-orbital active space. Full optimization at $6-31 \mathrm{G}^{*}$ not attemped, as this led only to a further $<0.1 \mathrm{kcal} \mathrm{mol}^{-1}$ reduction in energy at the $4-31 \mathrm{G}$ level. ${ }^{e}$ Relative energies calculated with the CAS8 active space. ${ }^{f}$ State-averaged orbitals.
and $2^{1} \mathrm{~A}_{\mathrm{g}}$ subsequently decays to the ground state on the picosecond time scale because of the presence of a twisted conical intersection. ${ }^{6 \mathrm{~d}}$ This means that the concentration of excited butadiene will always be small, and collisions involving an excited molecule will be unlikely. Nevertheless, low dimerization yields were measured. ${ }^{\text {an,c,d }}$ Yields for structures identified from the reaction in solution ${ }^{9 \mathrm{c}}$ are given in Scheme 3.

Shortly after the photochemical dimerization of butadiene was documented, it was discovered that dienes expedited the nonradiative decay of aromatic molecules. ${ }^{13 \mathrm{a}}$ A reaction profile, shown in Figure 1, involving excimer intermediates ${ }^{3,13,14}$ was suggested, but neither excimer emission nor products were detected at the time. ${ }^{13,14 f}$ Extensive work has since been carried

[^3]out on the photochemical addition of butadiene to benzene, ${ }^{15}$ naphthalene, ${ }^{16}$ and anthracene ${ }^{17-19}$ (Scheme 4).

Products have been characterized, and the formation of strained trans-diene adducts taken to indicate that the reaction is concerted. ${ }^{17 \mathrm{a}, \mathrm{f}}$ (The product structure reflects the dominant ground state conformer of the diene, ${ }^{20}, 11 \mathrm{c}$ not the most stable

[^4]product). The nature of any intermediates such as excimers is unclear, ${ }^{14 \mathrm{f}}$ as product formation might account for deactivation. ${ }^{18 \mathrm{a}}$ However, excimer emission has been detected in the diene + anthracene system ${ }^{17 \mathrm{c}, \mathrm{d}}$ (following indirect kinetic evidence for diene + naphthalene ${ }^{14 c-e}$ ) and is well established in the dimerization of anthracene, ${ }^{21}$ the prototype for photochromic systems. ${ }^{2 \mathrm{a}}$

Previous theoretical work on the $[4+4]$ reaction has been focused on the function of the excimer. ${ }^{8,22,4}$ In this paper, we concentrate on understanding the nonradiative decay channels for the model system butadiene + butadiene which involve conical intersections. These are consistent with the concerted nature of the $[4+4]$ reaction in real systems ${ }^{15-19}$ and with the limited experimental evidence for the formation of mixtures of products in butadiene + butadiene itself. ${ }^{9}$

## Computational Details

The choice of active space is the most critical feature of a CASSCF calculation. For two planar butadiene molecules which are wellseparated, the active space should consist of eight valence $\pi$ orbitals for covalent excited state calculations. However, the formation of a new $\sigma$-bond will lead to a pair of redundant active orbitals (one doubly occupied and the other unoccupied) and poor CASSCF convergence.

[^5]

Figure 1. The conventional model potential energy curves for photochemical $2+2$ or $4+4$ cycloaddition, based on calculations for $\mathrm{H}_{4}{ }^{3}$ This symmetric cut shows an excimer minimum on the singly excited state S (which fluorescences in certain systems) and the pericyclic minimum $\mathbf{A}$ which results from an avoided crossing of the ground state G and doubly excited state D .


Figure 2. The prototype $S_{1} / S_{0}$ conical intersection in but-1-ene. In a simple VB model, the exchange integrals $K_{i j}$ balance and the total exchange is zero at this geometry.
Redundant orbitals are therefore removed from the active space during geometry optimization but put back at the end for a final energy calculation.

MMVB ${ }^{22}$ structures were used as the starting point for CASSCF geometry optimizations. ${ }^{23}$ Guess orbitals were derived ${ }^{24}$ from stable

[^6]Table 3. Energies of the CASSCF/4-31G Antara Crossing Structures Shown in Figure 13

| structure | figure | active <br> space | $S_{0}, E_{\mathrm{h}}$ | $S_{1}, E_{\mathrm{h}}$ | $\Delta E\left(S_{1}-S_{0}\right)$ <br> $\mathrm{kcal} \mathrm{mol}^{-1}$ | $\Delta E$ on $S_{1},{ }^{a}$ <br> $\mathrm{kcal} \mathrm{mol}^{-1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{M}^{b} \mathbf{A}$ | $\mathbf{5 a}$ | 8 | -309.41640 | -309.35697 | 37.2 | 0.0 |
| $\mathrm{X}_{s}$ | $\mathbf{C 3 a}$ | 8 | $-309.24990^{c}$ | $-309.24621^{c}$ | 2.3 | +69.5 |
| $\mathrm{X} D_{2 d}$ | $\mathbf{1 3 b}$ | 8 | $-308.94006^{c}$ | $-308.93577^{c}$ | 2.7 | +264 |

${ }^{a}$ Energy on $S_{1}$ relative to pericyclic minimum $\mathbf{A} \cdot{ }^{b} \mathrm{M}=$ minimum; $\mathrm{X}=$ conical intersection. ${ }^{c}$ State-averaged orbitals.

UHF wave functions ${ }^{25}$ New $\sigma$ bonds were stretched to guarantee that the resulting active space could describe dissociation and that the final energetics would be comparable. The $4-31 \mathrm{G}$ basis used is sufficient to describe the topology of the $S_{1}$ surface, although polarization functions and an adequate treatment of dynamic electron correlation with geometry reoptimization would also be necessary to compute barrier heights accurately. Recomputing energies at the 6-31G* level does not change the barrier heights by more than $\pm 2 \mathrm{kcal} \mathrm{mol}^{-1}$ unless the geometry is highly strained (e.g., 12 in Tables 1 and 2).

Saddle points were characterized by analytic second derivative calculations at the CASSCF/4-31G level of theory. ${ }^{26}$ At these points, the single negative direction of curvature corresponds to the reaction coordinate. At a point on a conical intersection ${ }^{7}$ there are two linearlyindependent nuclear coordinates-the nonadiabatic coupling and gradient difference vectors (branching space ${ }^{71}$ )-which lift the degeneracy. The remaining directions define a space (intersection space ${ }^{71}$ ) in which the two states remain degenerate. Minima in the intersection space are located using the algorithm described in ref 27. Although decay can take place at any intersection point in principle, the region of the minimum will be favored when excess energy can be dissipated to the surroundings in condensed phases or when the system has very low excess energy such as cold jets. In these situations, the plane formed by the nonadiabatic coupling and gradient difference vectors will be the one in which initial motion on $S_{0}$ will take place. Furthermore, the gradient difference vector on $S_{1}$ will point along the reaction path.

## Results and Discussion

A VB Model for the Crossing Geometries for Butadiene + Butadiene. A large body of experience suggests that conical intersections-funnels ${ }^{3 \mathrm{aa}, 28}$ at which decay can be fully efficient ${ }^{7 \mathrm{~m}}$-are most likely to occur at tetraradicaloid geometries. ${ }^{29}$ A simple VB model ${ }^{66,29}$ can be used to predict the existence of such intersections. The problem is then to determine which crossing geometries will be at low energy and accessible for a particular system. We now briefly outline some possibilities that can be predicted a priori.

According to a simple VB model, ${ }^{6 c, 29}$ a tetraradicaloid geometry corresponds to a conical intersection if the exchange integrals between the four different radicaloid centers balance. The condition for the exchange integrals $\left(K_{i j}\right)$ is as follows:

$$
K_{12}+K_{34}=K_{14}+K_{23}=K_{13}+K_{24}
$$

The $K_{i j}$ depend mainly on the overlap of the orbitals on sites $i$ and $j$ via the usual expression for the exchange integral in Heitler-London VB theory.

$$
K_{i j}=\langle i j| \frac{1}{r_{12}}|j i\rangle+2 S_{i j}\langle i| h|j\rangle
$$

where $\langle i j| 1 / r_{12}|j i\rangle$ is the exchange repulsion, $\langle i| h|j\rangle$ is the nuclear electron attraction integral, and $S_{i j}$ is the overlap integral. Figure

[^7]Scheme 5


Scheme 6


2 illustrates this situation for the but-1-ene ${ }^{6 c}$ conical intersection. The first exchange equality is clearly satisfied by symmetry. The second can be fulfilled for certain values of the bond lengths and $\mathrm{C}-\mathrm{C}-\mathrm{C}$ angle. In butadiene + butadiene, the crossing $\mathbf{E}$ (Scheme 2) is of the but-1-ene type: one new $\sigma$ bond has already been formed, and the four radical centers are drawn three from one butadiene, and one from another. Another crossing-C in Scheme 2-corresponds to the $[2+2]$ addition of one double bond from each butadiene. ${ }^{6 \mathrm{a}, \mathrm{b}}$ This crossing can be thought of as a balance of exchange integrals brought about by the close approach of two radical centers across the diagonal of the rhombus.

Both $\mathbf{C}$ and $\mathbf{E}$ correspond to crossing geometries identified in the $\mathrm{H}_{4}$ system by Michl. ${ }^{3}$ Other tetraradicaloid geometries can be predicted, but these are disfavored in practice by framework strain or nonbonded repulsions. ${ }^{29}$ One such geometry is shown in Scheme 5 for trans + cis butadiene, in which three close radical centers (labeled 1,2, and 3) are distant from a fourth (labeled 4).

This is the $-(\mathrm{CH})_{3}-$ kink feature now recognized in polyenes ${ }^{6 \mathrm{~h}}$ and aromatic systems benzene ${ }^{6 \mathrm{i}}$ and styrene, ${ }^{6 \mathrm{j}}$ but, as we shall presently discuss, the framework distortion required to achieve this geometry in butadiene + butadiene places it unfavorably above the crossings $\mathbf{C}$ and $\mathbf{E}$ in energy.

Scheme 6 shows that the condition for zero total exchange can also be satisfied by the antara-antara approach of radical centers on the four terminal $-\mathrm{CH}_{2}$ groups.


Figure 3. Electronic states of butadiene + butadiene along the cis + cis suprafacial reaction coordinate. Energies calculated with MMVB at the pericyclic minimum on $S_{1}$ and the van der Waals minimum on $S_{0}$. At both geometries, $S_{1}$ corresponds to $\mathrm{D}_{\mathrm{A}}+\mathrm{D}_{\mathrm{B}}$ (i.e., $2^{1} \mathrm{~A}_{\mathrm{g}}$ excitation in both butadienes). At the asymptotic limit (not shown) the $D_{A}+G_{B}$ state lies below the $D_{A}+D_{B}$ if the geometry is relaxed. However (see Table 1) if the geometry is fixed at that shown in Figure $4 b$, the $\mathrm{D}_{\mathrm{A}}+\mathrm{D}_{\mathrm{B}}$ state is $S_{1}$.


Figure 4. Optimized geometries for $S_{1}$ minima located with CASSCF/4-31G: a is the pericyclic minimum, and $\mathbf{b}$ is the isolated $S_{1}$ butadiene planar minimum. ${ }^{31}$ Energies in Table 1.

However, this difficult approach is excessively high in energy because of nonbonded repulsions, and we shall show that the antara-antara approach for the $[4+4]$ reaction can be ruled out.

As we have suggested, the rhomboid supra-supra conical intersection involving the four terminal radical centers (Scheme 1b) does not exist.

The Pericyclic Minimum. Figure 3 shows the qualitative behavior of the covalent excited states along a cis-cis reaction path (computed with MMVB ${ }^{22}$ ). It is convenient to classify these states according to the excited and ground states of butadiene: $\mathrm{G}_{\mathrm{A}}, \mathrm{G}_{\mathrm{B}}$ (ground state $S_{0}$ ); $\mathrm{S}_{\mathrm{A}}, \mathrm{S}_{\mathrm{B}}$ (corresponding to the singly excited ${ }^{1} \mathrm{~B}_{\mathrm{u}}$ ionic state of trans butadiene); $\mathrm{D}_{\mathrm{A}}, \mathrm{D}_{\mathrm{B}}$ (corresponding to the covalent doubly excited state $2^{1} \mathrm{~A}_{\mathrm{g}}$ in trans butadiene); and $T_{A}, T_{B}$ (corresponding to the triplet states ${ }^{3} B_{u}$ for trans butadiene). The central feature is the pericyclic minimum $\mathbf{A}$ on $S_{1}$ which correlates with two covalent doubly excited states $D_{A}+D_{B}$ or two triplet states $T_{A}+T_{B}$ of two butadienes at large internuclear separation in accord with the accepted model. ${ }^{4}$ Interactions between the four radical centres are attractive in butadiene + butadiene $S_{1}$ but repulsive in $S_{0}$. The pericyclic minimum has been optimized at the CASSCF/ 4-31G level and is shown in Figure 4. The two butadiene molecules are $2.33 \AA$ apart, and the inversion of single and


Figure 5. Optimized geometries for $S_{1}$ transition structures located with CASSCF/4-31G. Energies in Table 1.


Figure 6. Optimized geometries for $S_{0} / S_{1}$ conical intersection minima located with CASSCF/4-31G. Energies in Table 1.

(a)
(b)

Figure 7. Optimized geometries for $S_{0} / S_{1}$ conical intersection minima located with CASSCF/4-31G: trans-cis isomers of the structures illustrated in Figure 6. Energies in Table 1.



Figure 8. A minimum on $S_{1}$ at an almost identical geometry to the transition structure $\mathbf{D}$ in Figure 5.
double bonds within each butadiene is less extreme than for the isolated molecule in the planar $2^{1} \mathrm{~A}_{\mathrm{g}}\left(\mathrm{D}_{\mathrm{A}}\right)$ minimum $(\mathbf{4 b}) .{ }^{6 \mathrm{~d}}$ A frequency calculation shows that $\mathbf{A}$ is a true minimum, with a binding energy of $16 \mathrm{kcal} \mathrm{mol}^{-1}$ (Table 1). The $S_{1}-S_{0}$ gap at this geometry is $37 \mathrm{kcal} \mathrm{mol}^{-1}$ (Table 1) which prohibits fast internal conversion.

Accurate calculations in the region where the two butadiene fragments are far apart are unreliable without the use of extended basis sets and are complicated by the existence of a state in which four unpaired spins are parallel: this lies below the $\mathrm{D}_{\mathrm{A}}$ $+\mathrm{G}_{\mathrm{B}}$ state at large interfragment distance (see Figure 3). For butadiene itself, the $S_{1}$ covalent $2^{1} \mathrm{~A}_{\mathrm{g}}$ state $\left(\mathrm{D}_{\mathrm{A}}\right)$ lies above the ${ }^{1} \mathrm{~B}_{\mathrm{u}}$ ionic state $\left(\mathrm{S}_{\mathrm{A}}\right)$ in the vertical excitation region. However, rapid internal conversion to $2^{1} \mathrm{~A}_{\mathrm{g}}$ has been shown to take place $\left(\sim 10 \mathrm{fs}^{12}\right)$. At a geometry of two $S_{0}$ butadienes which are 3.9 $\AA$ apart, the degenerate state corresponding to $D_{A}+G_{B}$ is already higher in energy than $D_{A}+D_{B}$. Thus the initial part of the reaction path lies outside the scope of the present


Figure 9. IRC in forward and reverse directions from the transition structure $\mathbf{B}$.


Figure 10. IRC in forward and reverse directions starting from the transition structure $\mathbf{D}$.
investigation and the nature of any excimer minimum for butadiene + butadiene remains an open question.

Reaction Paths ABC and ADE Lead to Low Energy Crossings. Two reaction paths on $S_{1}$ have been characterizedABC and ADE (Scheme 2)-which commence at the pericyclic minimum $\mathbf{A}$ and proceed via transition structures $\mathbf{B}$ and $\mathbf{D}$ (Figure 5) to (cis + cis) conical intersections $\mathbf{C}$ and $\mathbf{E}$ (Figure 6 ). Energies are given in Tables $1(4-31 G)$ and $2\left(6-31 G^{*}\right)$. Trans + cis conical isomers $\left(\mathbf{C}_{\mathbf{c t}}\right.$ and $\left.\mathbf{E}_{\mathbf{c t}}\right)$ intersections $\mathbf{C}$ and $\mathbf{E}$ have also been optimized (Figure 7), but no trans + cis isomer minimum corresponding to $\mathbf{A}$ could be located.

The transition state $\mathbf{B}$ is calculated to be $15.3 \mathrm{kcal} \mathrm{mol}^{-1}$ above $\mathbf{A}$, less than $1 \mathrm{kcal} \mathrm{mol}^{-1}$ below the dissociation limit (Table 1). Figure 5 shows that the corresponding transition vector is dominated by the formation of a new $\sigma$ bond between the shaded centers. An IRC calculation from this point (Figure 9) demonstrates that this coordinate leads either back to $\mathbf{A}$ or forwards to a point on the $S_{1} / S_{0}$ crossing $\mathbf{C}$.

The barrier height for coordinate ADE is $9.4 \mathrm{kcal} \mathrm{mol}^{-1}(\sim 6$ $\mathrm{kcal} \mathrm{mol}^{-1}$ lower than that for $\mathbf{A B C}$ ). The transition vector is

C




Figure 11. 11. The derivative coupling and gradient difference vectorsthose which lift the degeneracy-computed with CASSCF/4-31G at the $\mathbf{C}$ and $\mathbf{E}$ crossings.


Figure 12. 12. A suprafacial trans-cis conical intersection optimized with CASSCF/4-31G, at higher energy (Table 1) than the crossings shown in Figures 6 and 7.
again dominated by the formation of a new $\sigma$-bond (Figure 5). However the transition state region is very flat. There is an additional distinct minimum (Figure 8) at a geometry close to $\mathbf{D}$, with a slightly shorter $\sigma$-bond which has the same energy as the TS D. The fact that force constants are quite different at an almost identical geometry implies that the $S_{1}$ surface in this region is highly anharmonic. Quadratic force constantsincluding the transition vector-are of limited value. Thus an IRC calculations from this point is only partially successful (Figure 10). The reverse coordinate leads to the pericyclic minimum A, the IRC in "forwards" direction illustrated in Figure 10 actually decayed back toward A.

The crossing geometry $\mathbf{C}$ concurs with the simple VB model outlined above. The shortest interfragment separation is the $2.15 \AA$ diagonal illustrated in Figure 6. One ethylene group from each butadiene participates; the other (bond length 1.35
$\AA$ ) is almost unchanged from the ground state value for the isolated molecule (Figure 14.b). This crossing is $8.8 \mathrm{kcal} \mathrm{mol}^{-1}$ above $\mathbf{A}$ (Table 1). The derivative coupling and gradient difference vectors at $\mathbf{C}$ are illustrated in Figure 11. These are the directions which lift the degeneracy, and, at low kinetic energies, initial motion on $S_{0}$ will take place in this plane. The two limiting motions (Figure 11) correspond to different ways of recoupling the four radical centers. ${ }^{30}$ One direction leads to the formation of two new $\sigma$-bonds and hence a four-membered ring. The other leads to the $\pi$-bonds being reformed and return to isolated reactants. The relative yields of these two products will be a function of the nuclear dynamics. ${ }^{31}$

The crossing $\mathbf{E}$ is $10.6 \mathrm{kcal} \mathrm{mol}^{-1}$ below $\mathbf{A}, 19.4 \mathrm{kcal} \mathrm{mol}^{-1}$ below $\mathbf{C}$, and the lowest energy point on $S_{1}$ overall (Table 1). Figure 6 shows that the closest approach on the opposite side of the new $\sigma$ bond is $2.18 \AA$. Comparison with the transition structure in Figure 5 shows that a $C_{2}$ axis of symmetry has been lost: at some point on ADE after $\mathbf{D}$, the reaction path must bifurcate to lead to two (equivalent) possible forms of $\mathbf{E}$.

The directions which lift the degeneracy at $\mathbf{E}$ are shown in Figure 11. One points toward the formation of a sevenmembered ring, the other is orthogonal, and combinations of the two could lead to six-, seven-, or eight-membered rings being formed initially. One of the surprises from the dynamics treatment of this problem ${ }^{32}$ (using MMVB ${ }^{22}$ to simulate the CASSCF potential) is that four-membered rings can also be formed after decay at this point. The existence of a conical intersection is therefore consistent with the experimental observation ${ }^{9}$ of a mixture of products, as indicated in Scheme 3. Note that alternate forms itermediate between $\mathbf{C}$ or $\mathbf{E}$ could also exist as precursors to a five-membered ring product. However these were not found in our work as minima on the conical intersection line. Rather, dynamics studies ${ }^{32}$ indicate that the five-membered ring product occurs from the conical intersection $\mathbf{E}$.
-(CH) $)_{3}$ - Kink and Antara-Antara Crossings. Figure 12 shows that, in agreement with the VB prediction, a conical intersection can be formed by two butadienes as a result of interactions among the central four CH units (the $-(\mathrm{CH})_{3}-$ kink conical intersection that exists in polyenes $\left.{ }^{6 \mathrm{~h}}\right)$. This crossing resembles $\mathbf{E}$, but with two new $\sigma$ bonds fully formed. There are two short cross-ring distances ( $1.95 \AA$ and $1.83 \AA$ ), and one more distant radical center (as indicated in Scheme 5). However, in order to achieve this geometry, the $\sigma$ framework



(a)

Figure 13. 13. Optimized geometries on the antara-antara $S_{1} / S_{0}$ conical intersection.

(a)
(b)

Figure 14. 14. Optimized geometries for $S_{0}$ minima located with CASSCF/4-31G: cyclooctatetraene $\mathbf{a}$ and isolated planar butadiene $\mathbf{b}$. Energies in Table 1.


Figure 15. 15. Optimized geometries for $S_{0}$ minima located with CASSCF/4-31G: 1,3 divinylcyclobutane. Energies in Table 1.
has been strained considerably (e.g., bond lengths $1.60 \AA$, and angles $\ll 120^{\circ}$ ). This crossing is consequently over 50 kcal $\mathrm{mol}^{-1}$ higher in energy than the pericyclic minimum (Table 1), although this high value is due in part to the lack of polarization functions in the $4-31 \mathrm{G}$ basis set used. Calculations at the $6-31 \mathrm{G}^{*}$ level (Table 2) reduce the barrier by $\sim 10 \mathrm{kcal} \mathrm{mol}^{-1}$.

The conical intersection structure shown in Figure 12 is included because this type of interaction may be important in the addition of dienes to aromatic systems, ${ }^{15-19}$ in which the necessary distortions may be better accommodated by the $\sigma$ framework than in butadiene. Scheme $4 \mathrm{a}, \mathrm{b}$ show that the products of diene + aromatic photodimerization reflect the initial equilibrium of trans:cis isomers in the diene ground state, ${ }^{20}$ which for butadiene is $\sim 96: 4 .{ }^{20 b}, 11 \mathrm{c}$ Although crossings $\mathbf{C}$ and $\mathbf{E}$ have been found to occur at trans + cis geometries $\left(\mathbf{C}_{\mathbf{c t}}\right.$ and
$\mathbf{E}_{\mathbf{c t}}$ Figure 7) it may be that the Figure 12 crossing becomes competitive, as the radical centers can be efficiently delocalized.

Two points on a crossing located on the antara-antara path are illustrated in Figure 13. Table 3 shows that the large nonbonded repulsions place the lowest point on this crossing (Figure 13a, $C_{s}$ geometry) at $\sim 70 \mathrm{kcal} \mathrm{mol}^{-1}$ above the pericyclic minimum $\mathbf{A}$. This crossing is therefore unfavorable both energetically and entropically and can be disregarded unless a suitably constrained precursor geometry can be devised.

Products. Figures 14 and 15 illustrate two representative products on $S_{0}$. Cyclooctadiene (14a) is calculated to be 95 $\mathrm{kcal} \mathrm{mol}{ }^{-1}$ below the lowest point on the crossing $\mathbf{E}$ and 8.6 $\mathrm{kcal} \mathrm{mol}^{-1}$ below the dissociation limit on $S_{0}$ (Table 1). Kinetic energy must therefore be efficiently dissipated into the surrounding medium for this product to form. Figure 15 illustrates an isomer of 1,3 -divinylcyclobutane, which is $98 \mathrm{kcal} \mathrm{mol}^{-1}$ below the lowest point on the crossing $\mathbf{C}$ (Table 1). An interpolation along the path from $\mathbf{C}$ to 1,3-divinylcyclobutane is presented in Figure 16. The product is $7.8 \mathrm{kcal} \mathrm{mol}^{-1}$ endothermic (Table 1) and may dissociate if the excess kinetic energy is not removed.

## Conclusion

Two reaction paths ABC and ADE have been characterized on the $S_{1}$ excited state of the model butadiene + butadiene system. Both lead to a conical intersection, at which $S_{1} \rightarrow S_{0}$ decay can be fully efficient. One intersection $(\mathbf{E})$ is the lowestenergy point on $S_{1}$ overall.

The two conical intersections we have located for butadiene + butadiene can be related to intersections which have been characterized previously in conjugated hydrocarbons. This suggests that such intersections are quite general features of the excited states of these systems. At each crossing point ( $\mathbf{C}$ and $\mathbf{E}$ ), a mixture of products is predicted by examining the two coordinates which lift the degeneracy. The fact that there are two crossings in butadiene + butadiene-with different barrier heights-suggests that the product yields will be a complex function of experimental conditions. However, because of the efficient unimolecular decay of photoexcited butadiene, this study is intended principally to provide a general model for the $[4+4]$ photocycloaddition. We have shown that it is not straightforward to generalize the model based on a highlysymmetric cut through the $\mathrm{H}_{4}$ potential energy surfaces ${ }^{2 a, 3,4}$ in

6a


C
Figure 16. 16. A linear interpolation between the $\mathrm{CASSCF} / 4-31 \mathrm{G}$ crossing geometry $\mathbf{6} \mathbf{a}$ and the four-membered ring product 15.
this case. The reason for this is that, although the conical intersections we have characterized for butadiene + butadiene can be related to features of the $\mathrm{H}_{4}$ system, ${ }^{3}$ reorganization of all eight $\pi$ electrons is crucial to produce those calculated to be favorable energetically (i.e., it is not always the same four electrons that are important). However, in contrast to the [2+ 2] addition ${ }^{3,6 a, b}$ and in agreement with the simple model, ${ }^{29}$ the pericyclic minimum is found to be a true minimum for butadiene + butadiene. The $S_{0}-S_{1}$ energy gap of $\sim 37 \mathrm{kcal} \mathrm{mol}^{-1}$ prohibits decay at this point.

Our target in this work has been to document the excited

[^8]state surface topology for a model $4+4$ cycloaddition. One may infer (from the geometry of the conical intersection and the nature of the gradient difference and derivative coupling vectors) that certain products may originate following decay at the conical intersections we have identified. However, the issue of product formation can only be answered as a result of dynamics studies ${ }^{32}$ which are required to fully understand the possible ways of recoupling the radical centers that occurs after decay at the conical intersection.

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[^0]:    $\dagger$ Department of Chemistry, King's College.
    § Dipartimento di Chimica ‘G. Ciamician’ dell’ Universitá di Bologna.
    ${ }^{\text {* }}$ Permanent address, Departament de Qumica Fisica, Universitat de Barcelona, Marti i Franques, 108028 Barcelona.
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